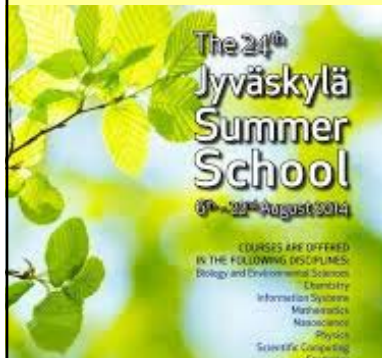


Isvector Nuclear Excitations and Isospin Symmetry



-the 4th lecture-

SS Jyvaskyla
August 06-12, 2014

Yoshitaka Fujita
Osaka University

Nucleon & Coin



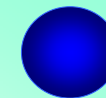
= Coin

back

front



proton



neutron

= Nucleon

similar mass
nearly the same interaction

$T_z = -1/2$

$T_z = 1/2$

isospin $T=1/2$

***Isvector Excitations & Isospin Symmetry

The ideas of "Isospin" and "Isospin Symmetry"
are important in IV excitations
that include τ operator.
(e.g., IV-E1(GDR), Gamow Teller)

Symmetry Natures of Strong and EM Interactions

Strong interaction: between
proton-proton, proton-neutron, neutron-neutron

IV int.

IV + IS int.

IV int.

EM interaction: between
proton-proton, proton-neutron, neutron-neutron

Yes

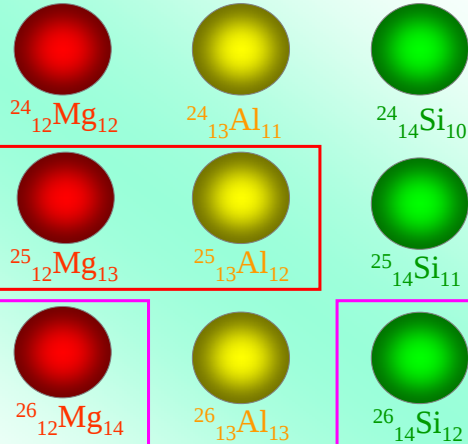
No

No

Isospin symmetry in simple terms

in their structure
Which are
the same?

Hint:
strong (nuclear)
interaction is
responsible for
the main part of
nuclear structure



**Fermi & Gamow-Teller Excitations

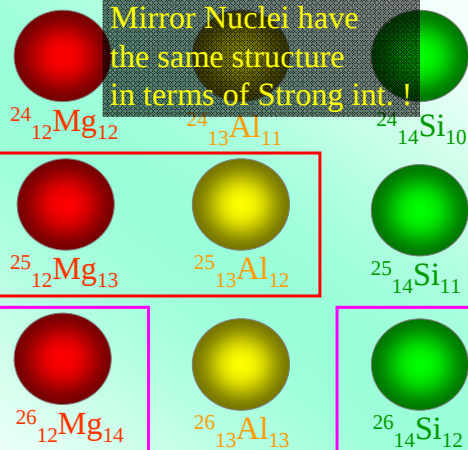
**Simplest nuclear excitations
that include τ operator.

**Especially, GT includes both τ and σ
that are unique in nuclei.

Isospin in simple terms

in their structure
Which are
the same?

Hint:
Strong (Nuclear)
interaction is
responsible for
the main part of
Nuclear structure



Vibration Modes in Nuclei (Schematic)

	Electric Mode ($\Delta S=0$)		Magnetic Mode ($\Delta S=1$)	
	IS ($\Delta T=0$)	IV ($\Delta T=1$)	IS ($\Delta T=0$)	IV ($\Delta T=1$)
L=0				
L=1				
L=2				
L=3				

Fermi mode
(τ)

Vibration Modes in Nuclei (Schematic)

	Electric Mode ($\Delta S=0$)		Magnetic Mode ($\Delta S=1$)	
	IS ($\Delta T=0$)	IV ($\Delta T=1$)	IS ($\Delta T=0$)	IV ($\Delta T=1$)
$L=0$				

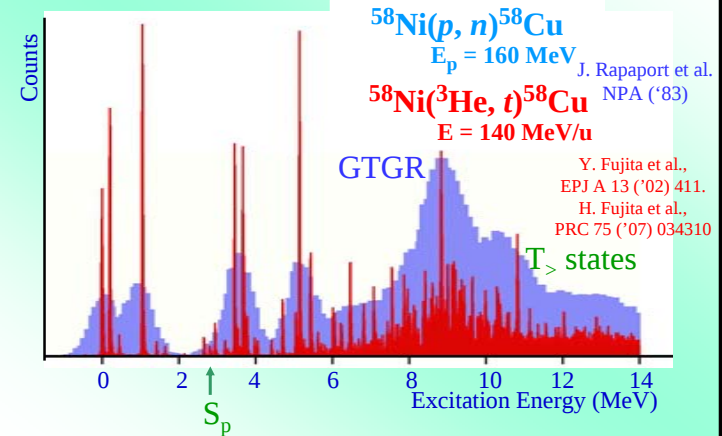
Gamow-Teller mode ($\sigma\tau$)

Fermi & GT transitions have $\Delta L=0$ character
 $\rightarrow$ therefore simple (no change in the radial w.f.)

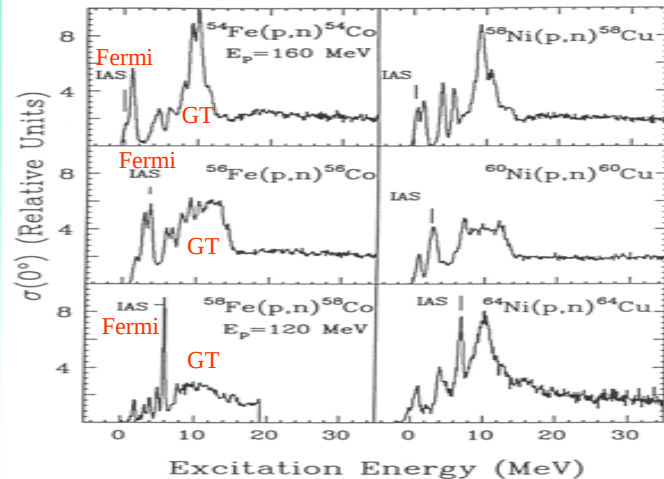
Smaller L transitions are favored in β and γ decays
 *** β and γ cannot carry a large L transfer

Fermi and GT transitions (weak processes) play important roles in the **Universe** ! (Astrophysics)

Comparison of (p, n) and $(^3\text{He}, t)$ 0° spectra



(p, n) spectra for Fe and Ni Isotopes



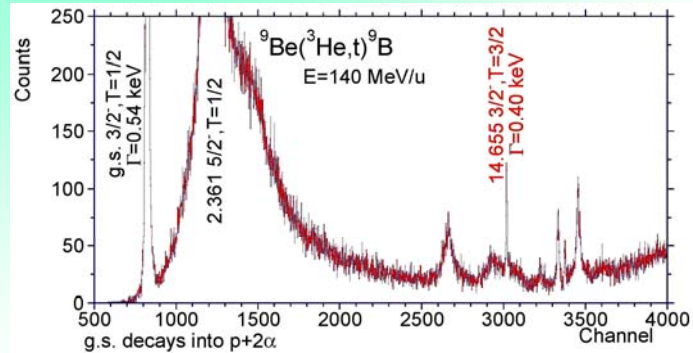
**Properties of GT transitions

Caused by the $\sigma\tau$ operator : a simple operator !

- 1) $|i\rangle$ and $|f\rangle$ states should have similar spatial shapes.
 - there is no space-type operator -
- 2) σ operator: states with $j_>$ and $j_<$ configurations are connected. (ex. $j_> = f_{7/2}$ and $j_< = f_{5/2}$)
- 3) τ operator: isospin quantum number T plays an important role. (isospin selection rule)

$\rightarrow$ GT transitions in each nucleus are **UNIQUE** !
 (reflecting nuclear structure of each nucleus)

${}^9\text{Be}({}^3\text{He}, t){}^9\text{B}$ spectrum (III)



14.7 MeV $T=3/2$ state is very weak!

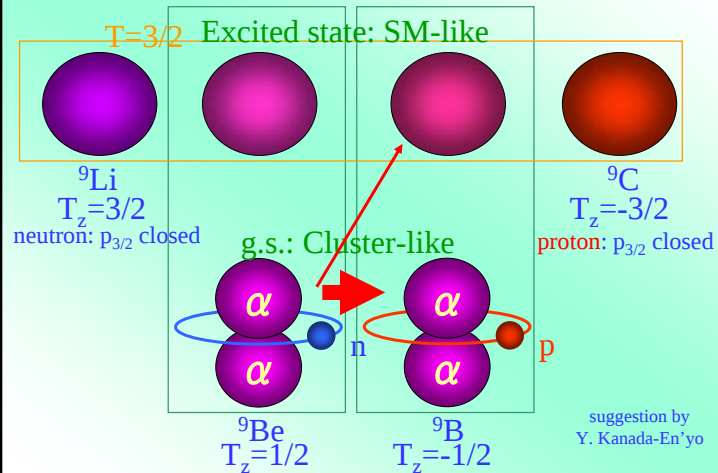
Strength ratio of g.s. & 14.7 MeV $3/2^-$ states: 140:1

***Isospin Symmetry

an important idea to see the connection of
decays and excitations caused
by Strong, EM and Weak interactions !

There are many cases that the "operators" are the same
in transitions caused by "strong," "EM" and "weak" int.

Shell Structure and Cluster Structure



$T=1/2$ Isospin Symmetry

Koelner Dom
in Germany
(157m high)



Nucleon & Coin



= Coin

back

face



proton



neutron

= Nucleon

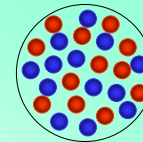
similar mass
nearly the same interaction

$T_z = -1/2$

$T_z = 1/2$

isospin $T = 1/2$

Isospin of a Nucleus



$A_N Z$

$$T_z = (1/2)N + (-1/2)Z$$

*z-component: conserved

The size of a vector should be larger than its z-component!

$$T = \text{or} > |T_z|$$

ex. ^{27}Al ($Z=13, N=14$) : $T_z = +1/2$, $T = 1/2, 3/2, \dots$

^{27}Si ($Z=14, N=13$) : $T_z = -1/2$, $T = 1/2, 3/2, \dots$

**only Z and N numbers are reversed !

Isospin Analogous Structure is expected !

Nucleus & Coin



= Coin

back

front



= Nuclei

$^{27}_{13}\text{Al}_{14}$

$^{27}_{14}\text{Si}_{13}$

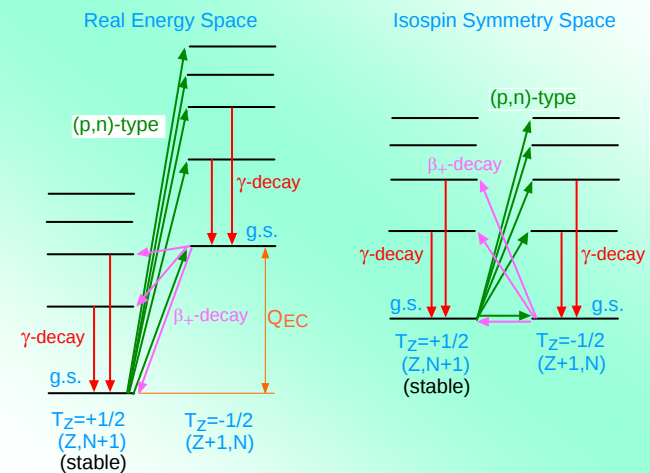
$$T_z = (1/2)N + (-1/2)Z$$

$T_z = 1/2$

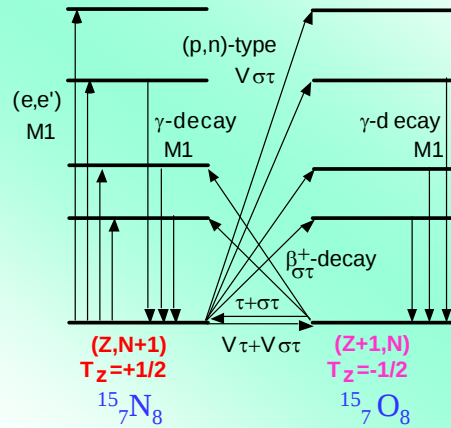
$T_z = -1/2$

isospin $T = 1/2, 3/2, \dots$

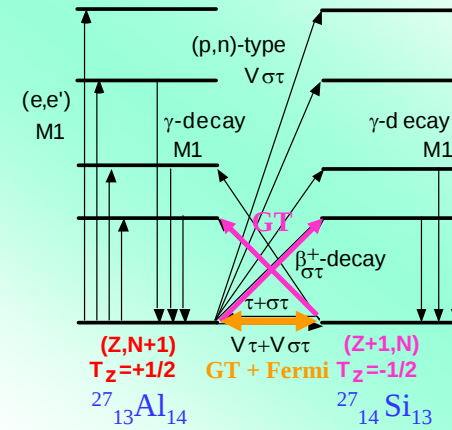
Analogous Structures, Transitions in $T=1/2$ System



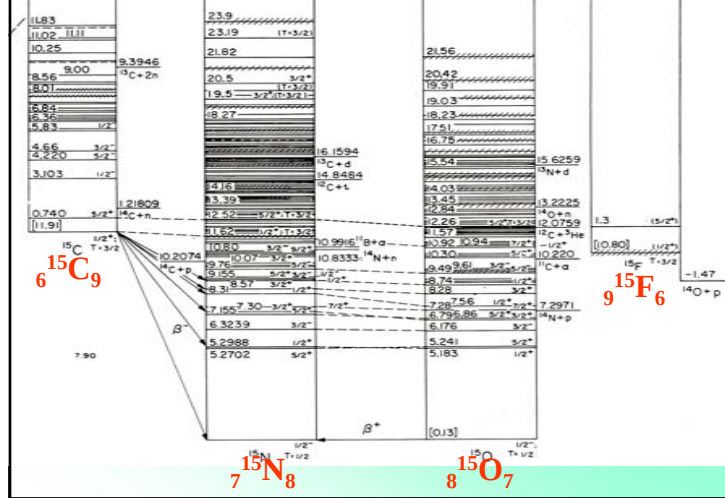
$T = 1/2$ Mirror Nuclei : Structures



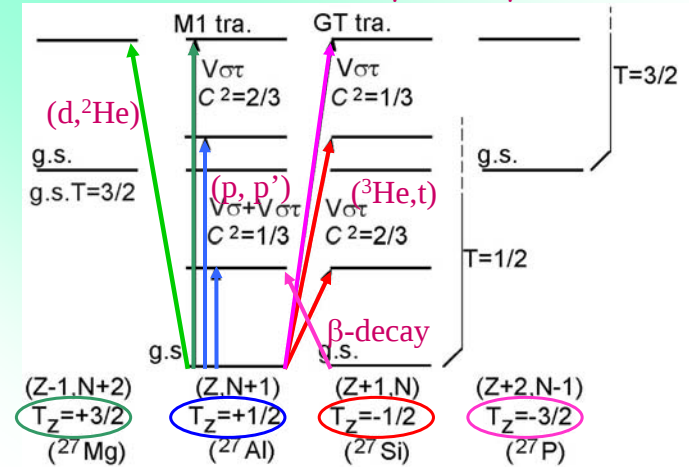
$T=1/2$ Mirror Nuclei : Structures & Transitions



$T=1/2, 3/2$ Isospin Symmetry for $A=15$ Nuclei



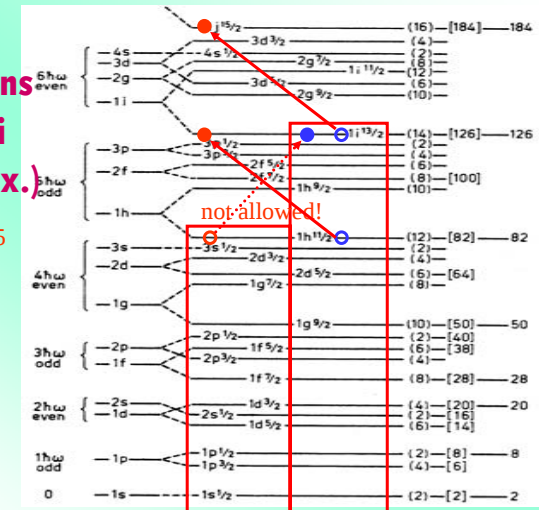
$T=1/2$ & $3/2$ Symmetry



**Transitions Starting from $T = 0$ Nuclei

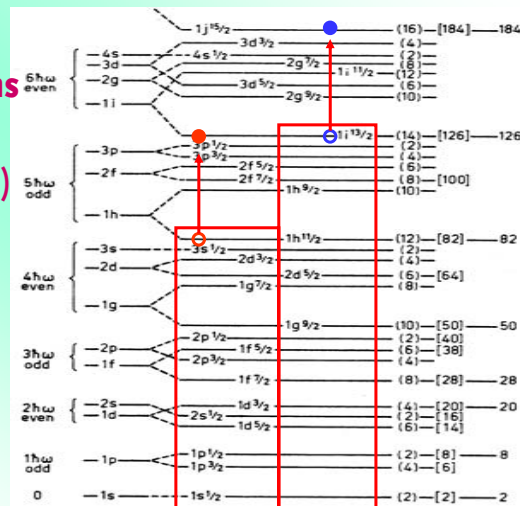
1p-1h Excitations in ^{208}Bi (charge ex.)

$Z=83, N=125$

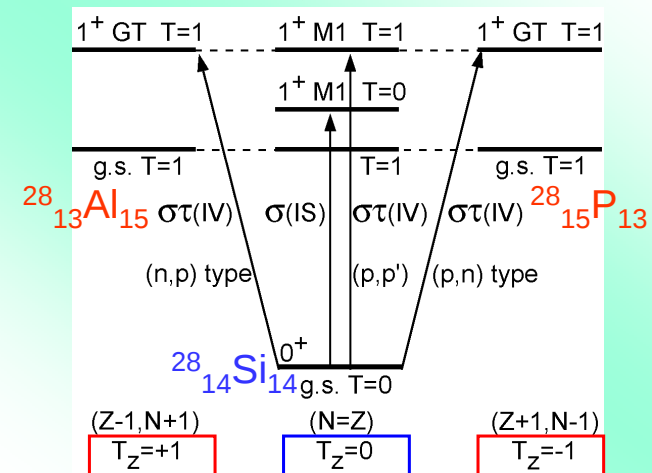


1p-1h Excitations in ^{208}Pb (inelastic)

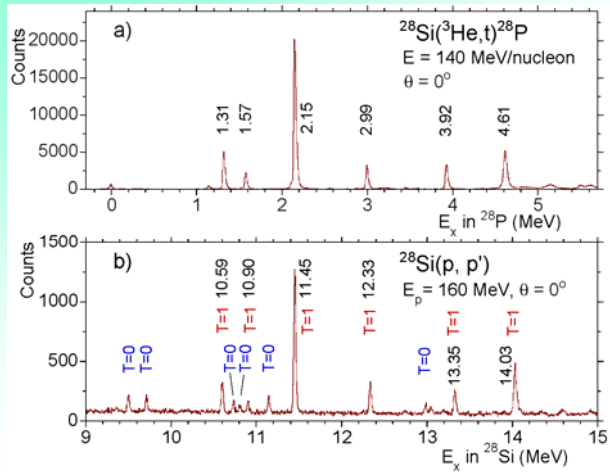
$Z=82, N=126$



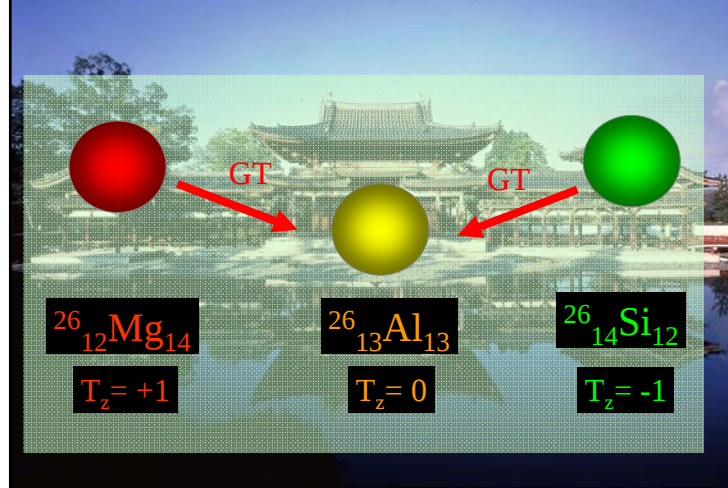
Isospin Symmetry in $T=0$ and 1 Nuclei



Spectra $^{28}\text{Si}(p, p')$ and $^{28}\text{Si}(^3\text{He}, t)^{28}\text{P}$

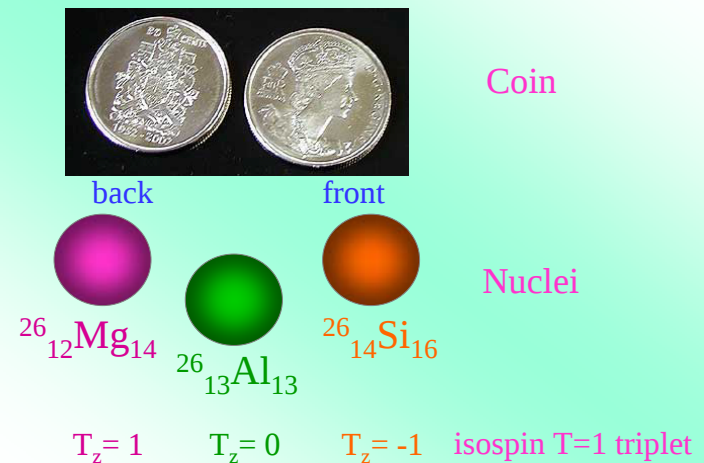


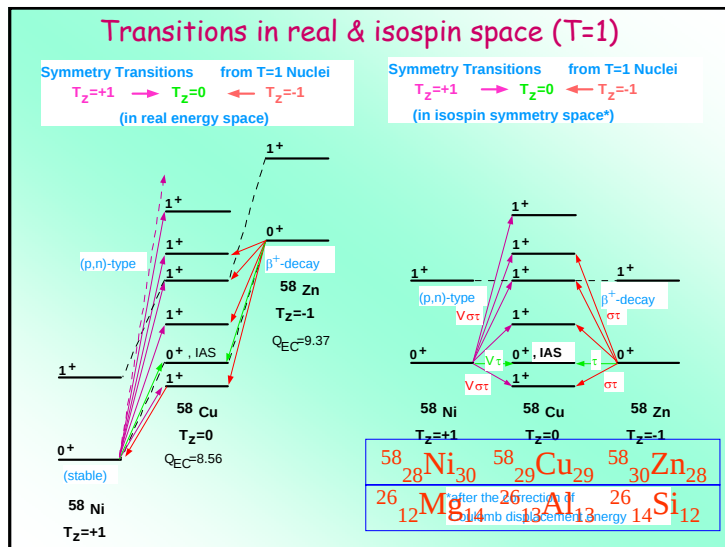
$T=1$ Isospin Symmetry



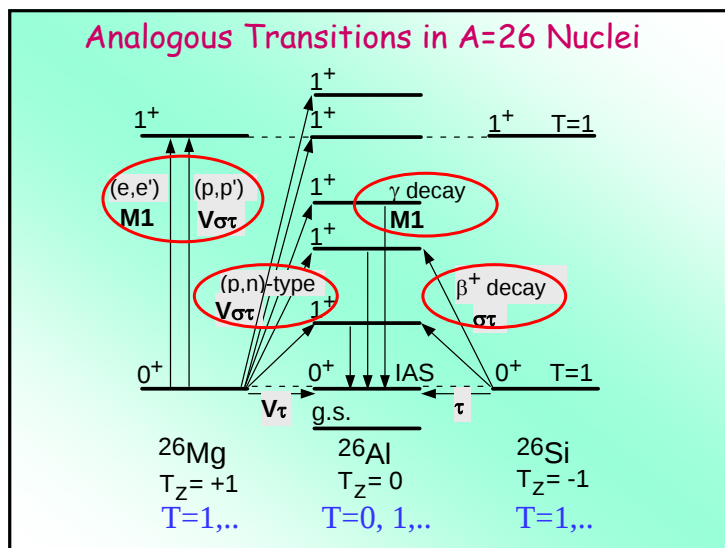
**Higher T Symmetry

Nuclei & Coin





***Derivation of B(GT) strength



Case 1 : γ -decay & β -decay

*both have very simple mechanism.
 (people even don't think of "mechanism !")

*Operators are relatively simple!

Weak : Gamow-Teller, Fermi

EM : E1, E2, ... M1, M2, ...

matrix element & $t_{1/2}$
 $(1 / t_{1/2}) = \text{Coup. Const.} \times \text{PhaseSpaceFac.}$
 $\times |\langle \mathbf{f} | \mathbf{Op} | \mathbf{i} \rangle|^2$

*if Op is specified, w.f.(=structures) are studied !
 (Op specification is not always easy!)

* highly *Ex* region cannot be reached !

Reduced transition strength $B(Op)$

A value proportional to (matrix element)²
 $|\langle f | Op | i \rangle|^2$
 is called "reduced transition strength"
 ex. $B(GT)$, $B(F)$, $B(M1)$, $B(E2)$,....

- *representing only the structure part
for a specific operator!
- *reaction mechanism part is removed!

$B(GT)$ derivation

★ β decay :fundamental, but E_x range :limited "Q-window limitation"

★ (p, n) reaction at intermediate energies ($E = 100$ -500 MeV)
 "proportionality" : $B(GT)$ and $\sigma(0^\circ)$

$$\sigma(0^\circ) = KN_{\sigma\tau} |J_{\sigma\tau}(0^\circ)|^2 B(GT)$$

⇒ Breakthrough against "Q-window limitation"
 but resolution : rather poor ($\Delta E = 200$ -400 keV)

★ $(^3\text{He}, t)$ reaction at intermediate energies ($E = 130$ -150 MeV/u)
 "high resolution" ($\Delta E < 50$ keV)

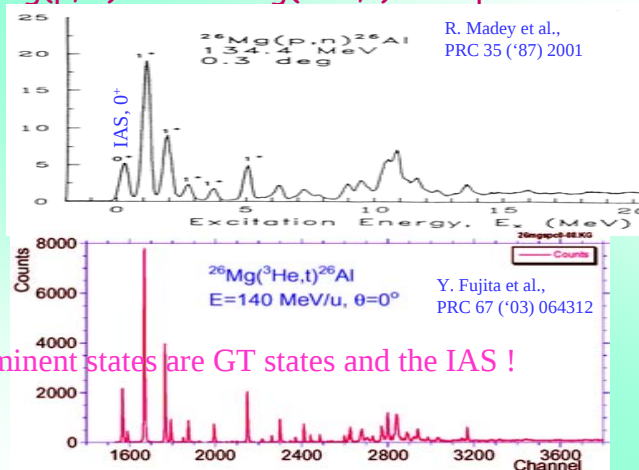
☆ magnetic spectrometer, matching techniques

"proportionality" : good ($B(GT) > 0.03$)

⇒ Breakthrough against "Energy resolution limitation"

⇒ Reliable $B(GT)$ values for individual transitions

$^{26}\text{Mg}(p, n)^{26}\text{Al}$ & $^{26}\text{Mg}(^3\text{He}, t)^{26}\text{Al}$ spectra



$B(GT)$ derivation

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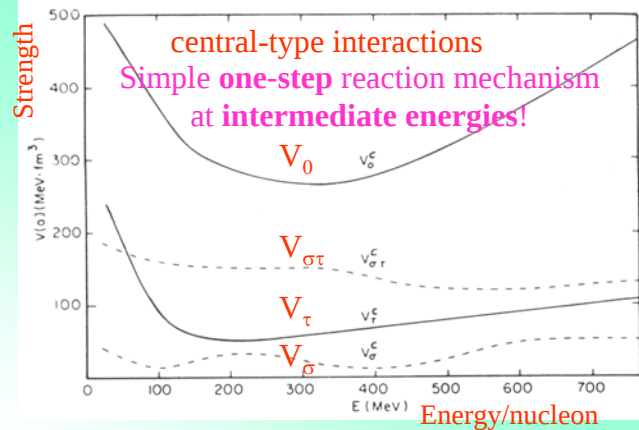
☆ magnetic spectrometer, matching techniques

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⇒ Reliable $B(GT)$ values for individual transitions

Nucleon-Nucleon Int. : E_{in} dependence at $q=0$



Love & Franey PRC 24 ('81) 1073

$B(GT)$ derivation

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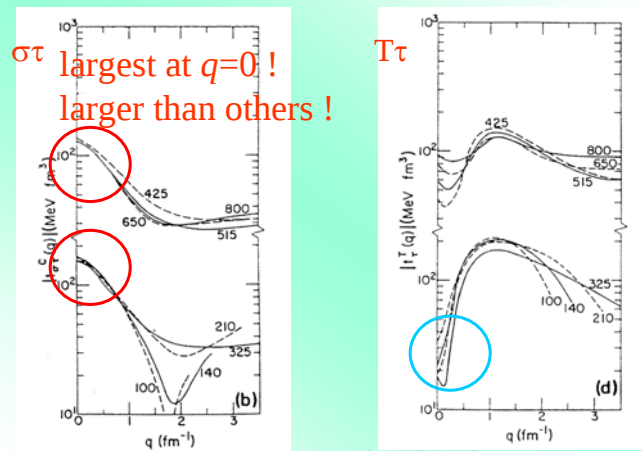
☆ magnetic spectrometer, matching techniques

"proportionality" : good ($B(GT) > 0.03$)

⇒ Breakthrough against "Energy resolution limitation"

⇒ Reliable $B(GT)$ values for individual transitions

N.-N. Int. : $\sigma\tau$ & Tensor- τ q -dependence



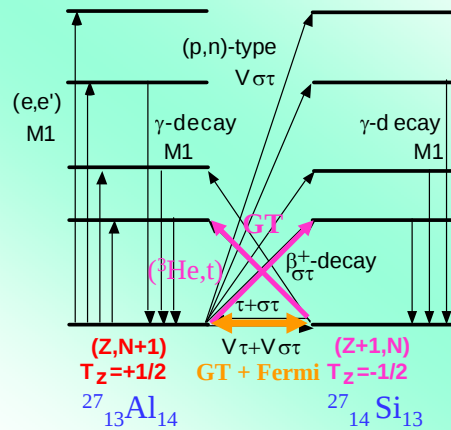
Love & Franey PRC 24 ('81) 1073

$T=1/2$ Isospin Symmetry

Koelner Dom
in Germany
(157m high)



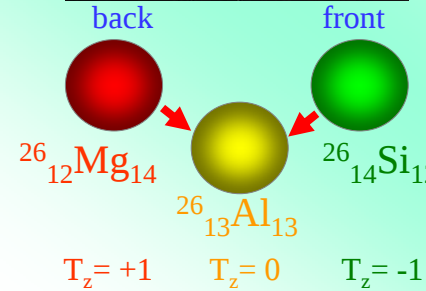
$T=1/2$ Mirror Nuclei : Structures & Transitions



Nuclei & Coin



= Coin



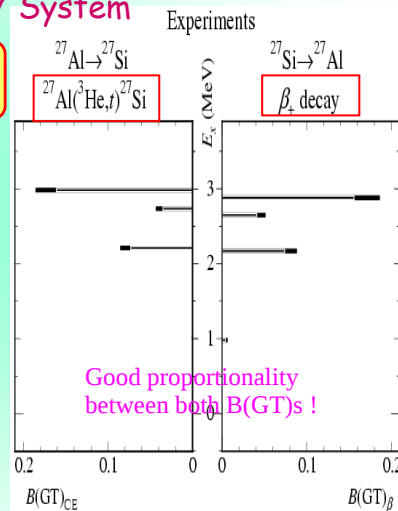
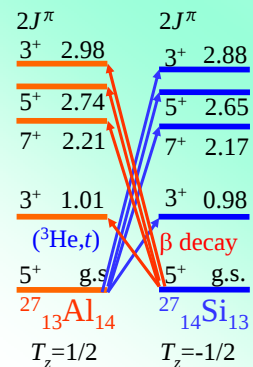
= Nuclei

Symmetry in
1) structure
2) transitions

isospin $T=1$ triplet

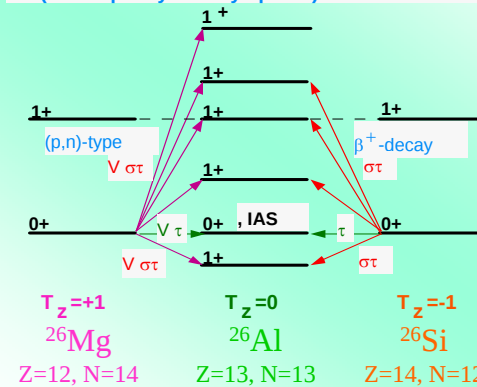
Symmetry in $A=27$ System

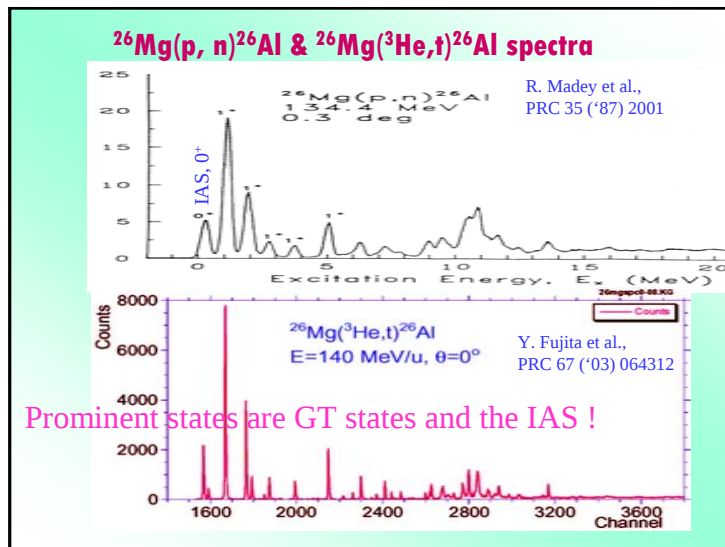
$$\frac{d\sigma}{d\Omega}(q \rightarrow 0) = KM |J|^2 B(GT)$$



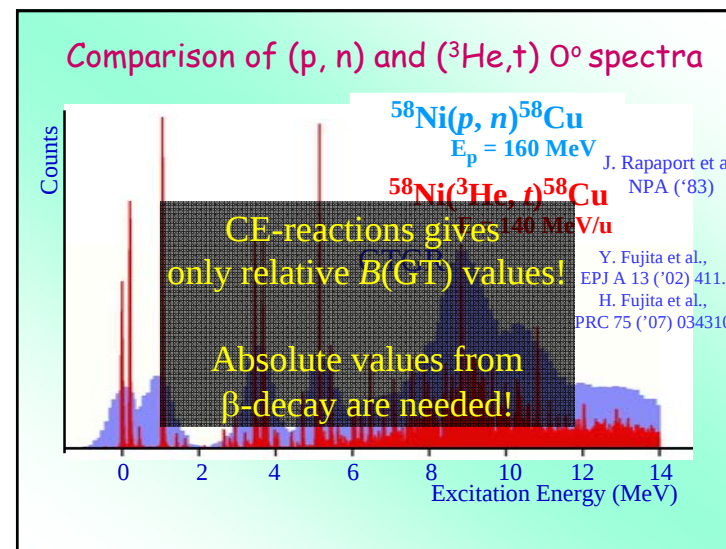
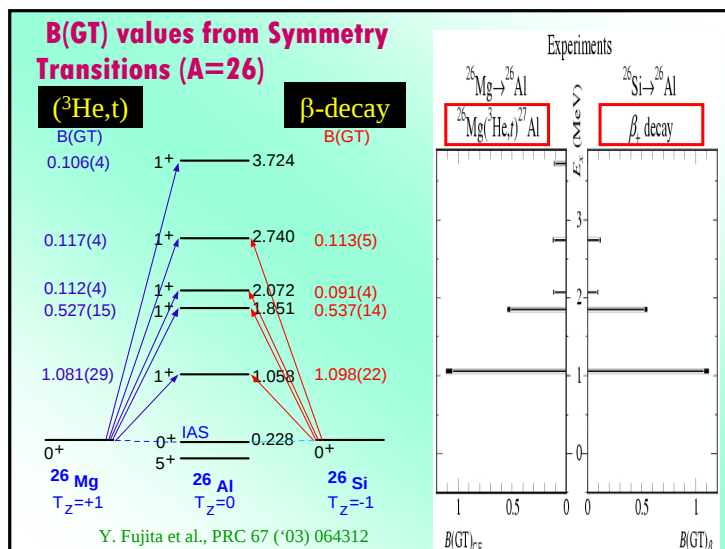
$T=1$ symmetry : Structures & Transitions

$T_z = +1 \rightarrow T_z = 0 \leftarrow T_z = -1$
(in isospin symmetry space*)





****CE reaction and β -decay: complementary tools**



β-decay & Nuclear Reaction

$$*\beta\text{-decay GT tra. rate} = \frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$$

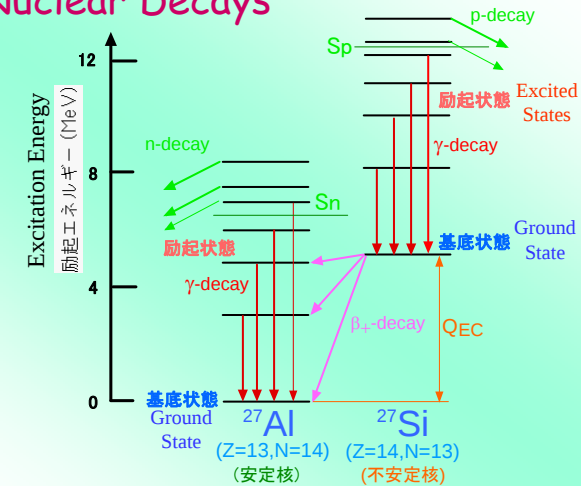
$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2$

*Nuclear (CE) reaction rate (cross-section)
 = reaction mechanism

$$\begin{matrix} \otimes \text{ operator} \\ \otimes \text{ structure} \end{matrix} = (\text{matrix element})^2$$

A simple reaction mechanism should be achieved !
 → we have to go to high incoming energy

Nuclear Decays



β-decay & Nuclear Reaction

$$*\beta\text{-decay GT tra. rate} = \frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$$

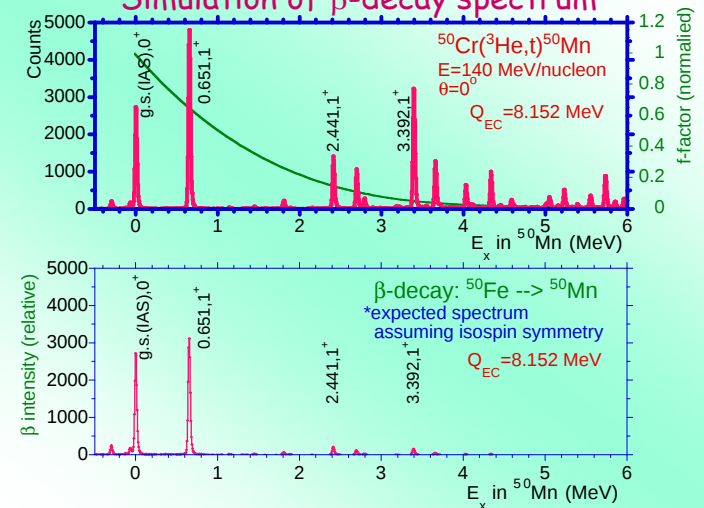
$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2$
 Study of Weak Response of Nuclei
 by means of

*Nuclear (CE) reaction rate (cross-section)
 = reaction mechanism
 using β-decay as a reference

$$\begin{matrix} \otimes \text{ operator} \\ \otimes \text{ structure} \end{matrix} = (\text{matrix element})^2$$

A simple reaction mechanism should be achieved !
 → we have to go to high incoming energy

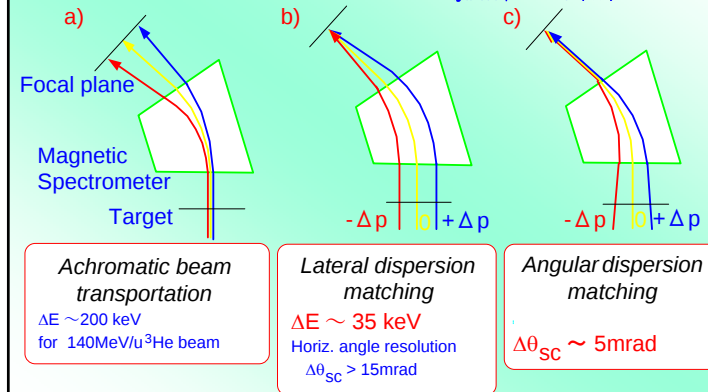
Simulation of β-decay spectrum



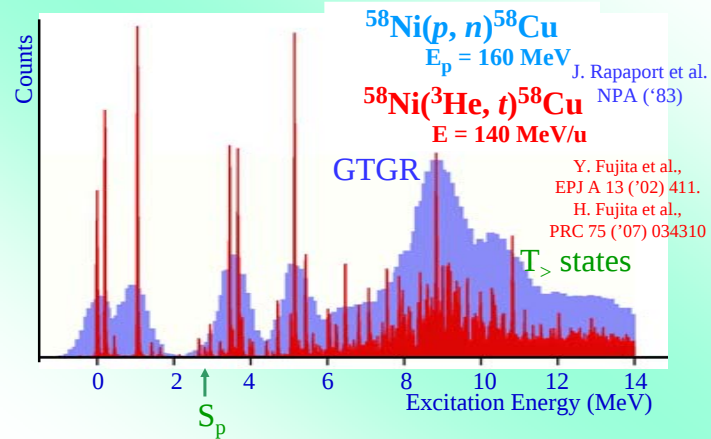
High Resolution Experiment

Matching Techniques

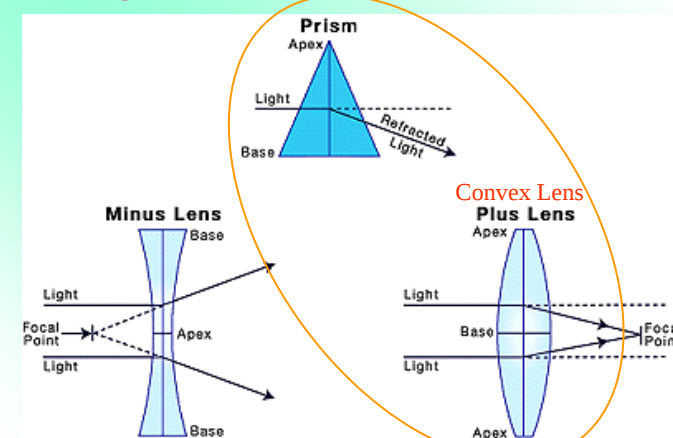
Y. Fujita et al., NIM. B 128 (1997) 274.
H. Fujita et al., NIM. A 484 (2002) 17.

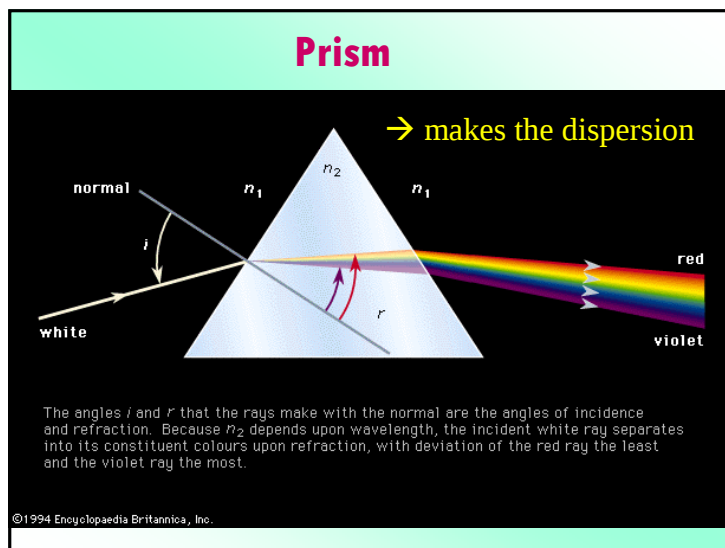
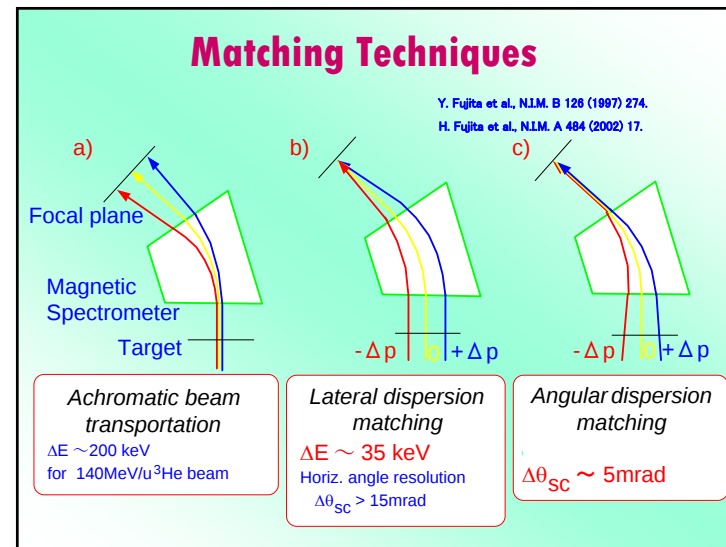
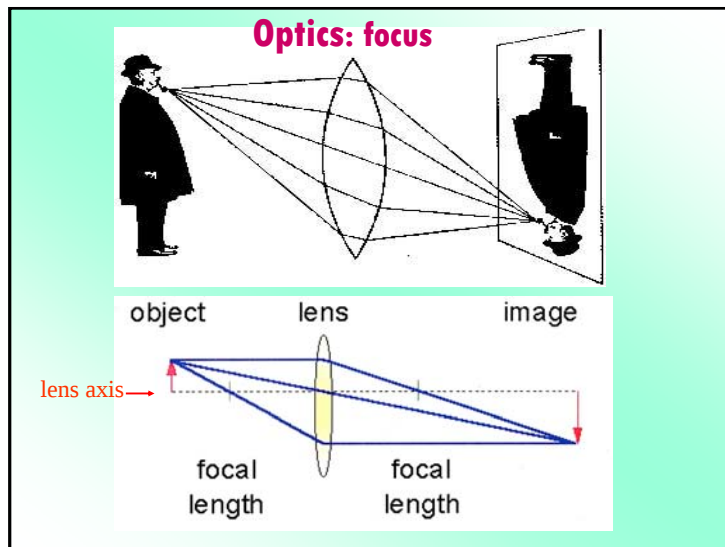


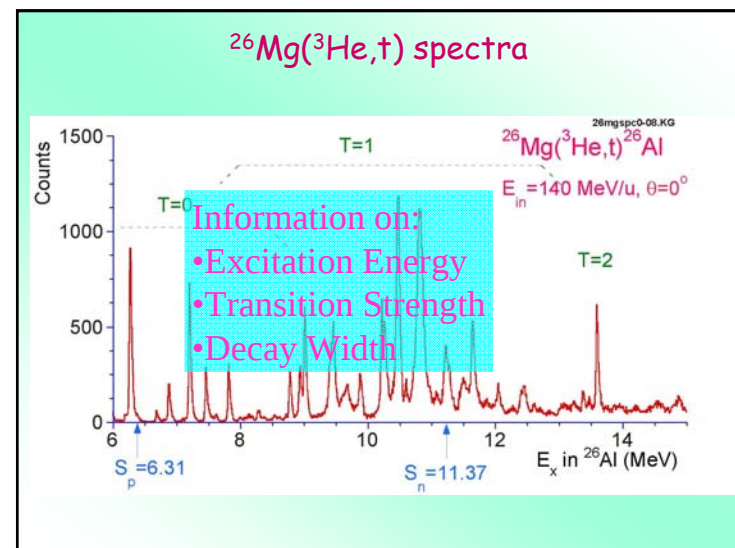
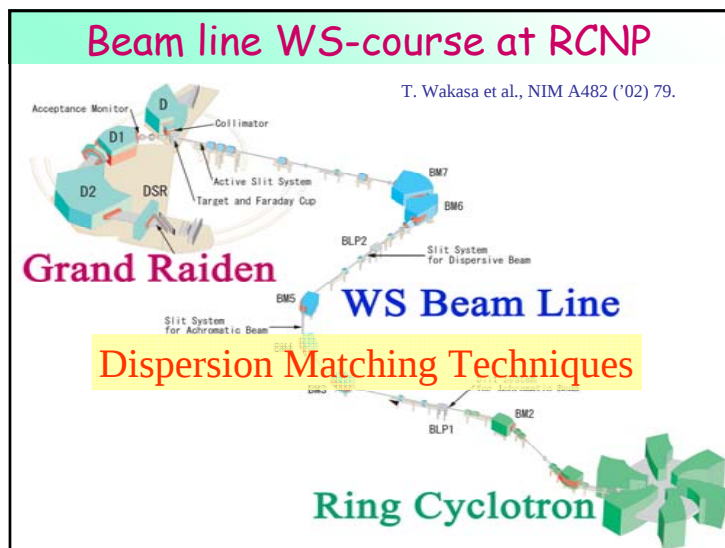
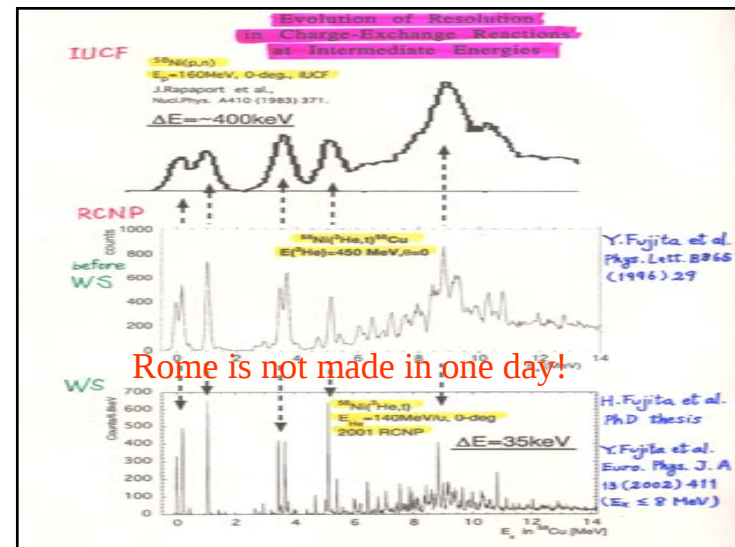
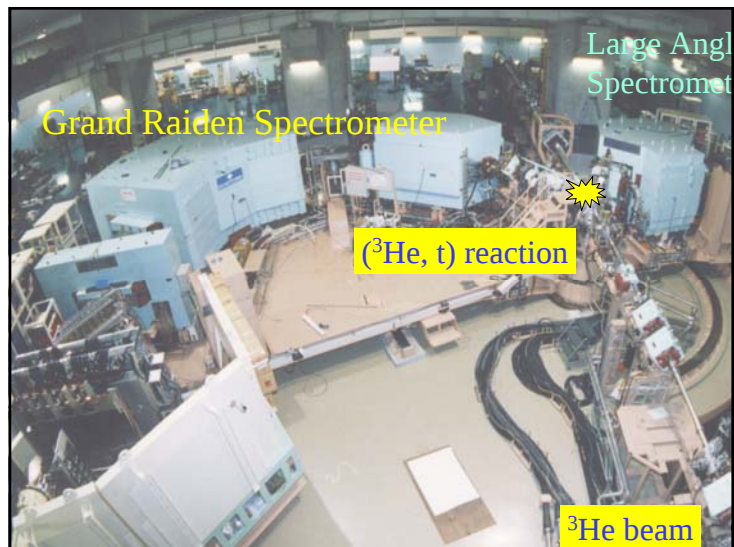
Comparison of (p, n) and (^3He , t) 0° spectra



Magnet = convex Lens + Prism

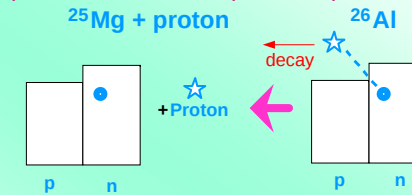






**T (Isospin) Selection Rules

Importance of Isospin : in p -decay of ^{26}Al



$$T_z : 1/2 + (-1/2) = 0$$

$$T : 1/2 \oplus 1/2 = 0 \text{ or } 1$$

$$3/2 \oplus 1/2 = 1 \text{ or } 2$$

S_p (p-sep. energy) in ^{26}Al : 6.31 MeV

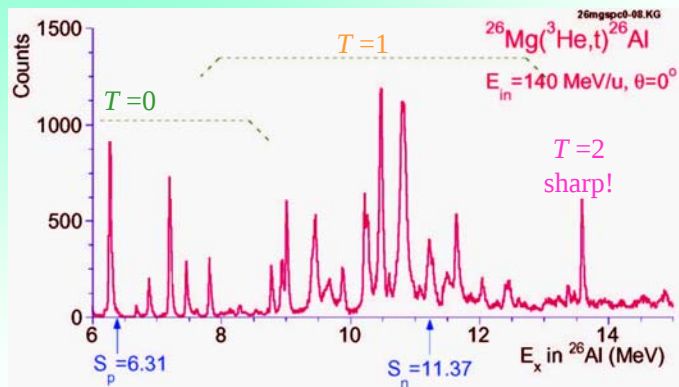
$T=3/2$ state in ^{25}Mg : $E_x > 7.79$ MeV

→ effective S_p in ^{26}Al

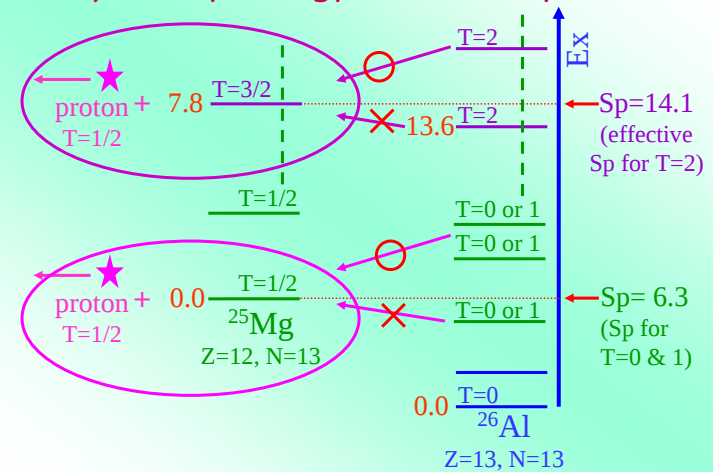
for $T = 0, 1$ states : $E_x = 6.31$ MeV

for $T = 2$ states : $E_x = 14.1$ MeV

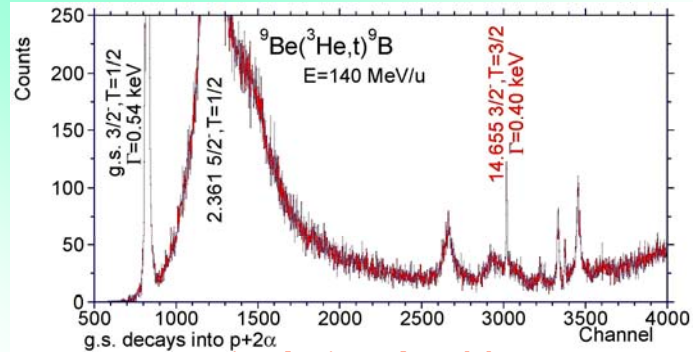
Higher- E_x region in ^{26}Al : $T=1$ & $T=2$ states



p -decay energy relationship

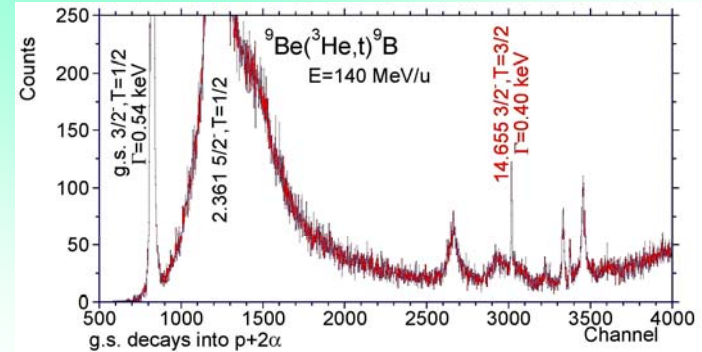


${}^9\text{Be}({}^3\text{He}, t){}^9\text{B}$ spectrum (II)



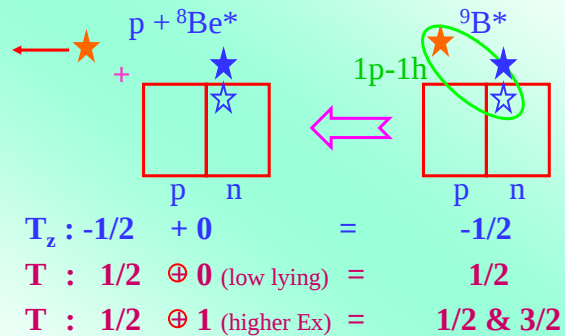
Isospin selection rule prohibits
proton decay of $T=3/2$ state!

${}^9\text{Be}({}^3\text{He}, t){}^9\text{B}$ spectrum (III)



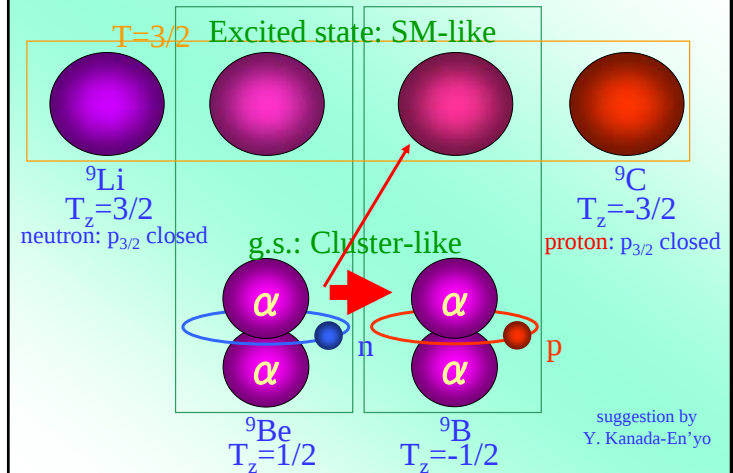
14.655 MeV $T=3/2$ state is very weak!

Isospin Selection Rule : in p -decay of ${}^9\text{B}$



* $T=1$ state in ${}^8\text{Be}$ is
only above
 $E_x=16.6$ MeV

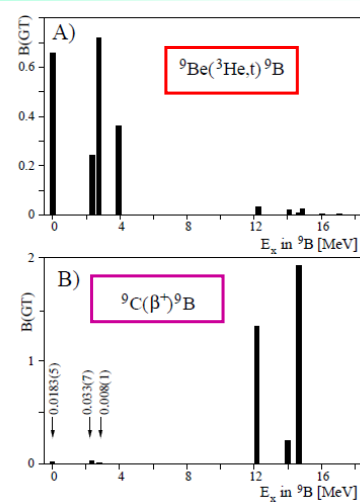
Shell Structure and Cluster Structure



β -decay and $(^3\text{He}, t)$ results

C. Scholl et al,
PRC 84, 014308 (2011)

L. Buchmann et al.,
PRC 63 (2001) 034303.
U.C. Bergmann et al.,
Nucl. Phys. A 692 (2001) 427.



Summary

Isospin quantum number: connects mass A nuclei
-unique in nuclei-

High Resolution: brings something new!
-one order difference makes the quality different-

Charge Symmetry / Charge Invariance

荷電対称性

Charge Symmetry
equality of
p-p & n-n
interactions

荷電不変性

Charge Invariance
equality of
p-p n-p n-n
interactions

